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Effects of Turbulence on the Transverse Position-Momentum Entanglement of Biphotons

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Motivations



- 1. Entanglement provides secure information transmission
- 2. High dimensional Hilbert space (qubits only 2D)
 - time-energy
 - transverse spatial coordinates of biphotons

□ orbital angular momentum (OAM)

- transverse linear position / momentum (quantum image)
- 3. Long distance communication in free space (turbulence effect)



Generation of Entangled Photons

Type-II Spontaneous Parametric Down-Conversion (SPDC)



Generation of Entangled Photons

Type-II Spontaneous Parametric Down-Conversion (SPDC)



We will restrict ourselves to degenerate and nearly collinear SPDC taking,

$$2\omega_s = 2\omega_i = \omega_p$$

$$2k_s \simeq 2k_i \simeq k_p$$

Theory of SPDC (Gaussian Approx. 1)



The biphoton state is given by

$$|\Psi
angle = \int d\vec{p_s} \ d\vec{p_i} \ \Phi(\vec{p_s}, \vec{p_i}) \ a_s^{\dagger}(\vec{p_s}) a_i^{\dagger}(\vec{p_i}) |0,0
angle$$

with the wave function in momentum space

$$\Phi(\vec{p}_s, \vec{p}_i) = N E_p(\vec{p}_s + \vec{p}_i) \operatorname{sinc}\left(\frac{\Delta_k L}{2}\right) \exp\left(i\frac{s_k L}{2}\right)$$

where N is a normalization factor,

 $\Delta_k = k_p(\vec{p_s} + \vec{p_i}) - k_s(\vec{p_s}) - k_i(\vec{p_i}) \text{ and } s_k = k_p(\vec{p_s} + \vec{p_i}) + k_s(\vec{p_s}) + k_i(\vec{p_i})$

 E_p is the transverse profile of the pump.

Theory of SPDC (Gaussian Approx. 1)



The biphoton state is given by

$$|\Psi
angle = \int d\vec{p_s} \ d\vec{p_i} \ \Phi(\vec{p_s}, \vec{p_i}) \ a_s^{\dagger}(\vec{p_s}) a_i^{\dagger}(\vec{p_i}) |0,0
angle$$

with the wave function in momentum space

$$\Phi(\vec{p}_s, \vec{p}_i) = N E_p(\vec{p}_s + \vec{p}_i) \operatorname{sinc}\left(\frac{\Delta_k L}{2}\right) \exp\left(i\frac{s_k L}{2}\right)$$
$$\to N \exp\left[-\frac{1}{2B}(\vec{p}_s + \vec{p}_i)^2\right] \exp\left[-\frac{1}{2A}(\vec{p}_s - \vec{p}_i)^2\right]$$

Theory of SPDC (Gaussian Approx. 1)



The biphoton state is given by

$$|\Psi\rangle = \int d\vec{p}_s \ d\vec{p}_i \ \Phi(\vec{p}_s, \vec{p}_i) \ a_s^{\dagger}(\vec{p}_s) a_i^{\dagger}(\vec{p}_i) |0,0\rangle$$

with the wave function in momentum space

$$\Phi(\vec{p}_s, \vec{p}_i) = NE_p(\vec{p}_s + \vec{p}_i) \operatorname{sinc}\left(\frac{\Delta_k L}{2}\right) \exp\left(i\frac{s_k L}{2}\right)$$
$$\to N \exp\left[-\frac{1}{2B}(\vec{p}_s + \vec{p}_i)^2\right] \exp\left[-\frac{1}{2A}(\vec{p}_s - \vec{p}_i)^2\right]$$
ace)
$$\Psi(\vec{x}_s, \vec{x}_i) \to N \exp\left[-\frac{B}{2}(\vec{x}_s - \vec{x}_i)^2\right] \exp\left[-\frac{A}{2}(\vec{x}_s + \vec{x}_i)^2\right]$$

(position space)

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OPTICS Theory of Turbulence



The propagated field in a turbulent medium is given by

[Ref: Milonni, AJP 67, 476(1999)]

$$\hat{E}^{(+)}(\vec{x},z) = e^{ikz} \int d\vec{x}' \ h(\vec{x},\vec{x}',z) e^{i\phi(\vec{x}')} \hat{E}^{(+)}(\vec{x}',0)$$

where

 $h(ec{x},ec{x}',z) = rac{k}{2\pi i z} \exp\left[rac{ik}{2z}(ec{x}-ec{x}')^2
ight]$ (paraxial limit)

The field at z = 0 is given by the Fourier transform of the annihilation operator

$$\hat{E}^{(+)}(\vec{x},0) \sim \int d\vec{q} \ \hat{a}(\vec{q}) \ e^{i\vec{q}\cdot\vec{x}}$$

OPTICS Theory of Turbulence



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[Ref: Milonni, AJP 67, 476(1999)]

$$\hat{E}^{(+)}(\vec{x},z) = e^{ikz} \int d\vec{x}' \ h(\vec{x},\vec{x}',z) e^{i\phi(\vec{x}')} \hat{E}^{(+)}(\vec{x}',0)$$

- 1. Turbulent medium is reflected from the statistical character of $\phi(\vec{x}')$.
- 2. The medium is replaced by a single "phase screen" accounting for all the phase fluctuation incurred in the propagation to *z*, i.e., $\phi(\vec{x}') = k \int_0^z n(\vec{x}', z') dz'$
- 3. Fluctuating phase: $e^{i[\phi(\vec{x}')-\phi(\vec{y}')]} = e^{-(1/2)D_s(|\vec{x}'-\vec{y}'|)}$ Phase structure function $D_s(|\vec{x}'-\vec{y}'|) = \alpha |\vec{x}'-\vec{y}'|^{5/3}$ (Kolmogorov)

Four-point correlation function

$$G(\vec{x}_{s}, \vec{y}_{s}; \vec{x}_{i}, \vec{y}_{i}) \equiv \langle \Psi | \overline{\hat{E}^{(-)}(\vec{y}_{i}, z_{i}) \hat{E}^{(-)}(\vec{y}_{s}, z_{s}) \hat{E}^{(+)}(\vec{x}_{s}, z_{s}) \hat{E}^{(+)}(\vec{x}_{i}, z_{i})} | \Psi \rangle$$

$$= \iiint d\vec{x}_{s}' d\vec{y}_{s}' d\vec{x}_{i}' d\vec{y}_{i}' T(\cdots) \times G_{0}(\vec{x}_{s}', \vec{y}_{s}'; \vec{x}_{i}', \vec{y}_{i}')$$

SPDC + **Turbulence**

Free space transfer function

$$T(\cdots) \equiv h(\vec{x}_s - \vec{x}'_s, z_s)h^*(\vec{y}_s - \vec{y}'_s, z_s)h(\vec{x}_i - \vec{x}'_i, z_i)h^*(\vec{y}_i - \vec{y}'_i, z_i)$$

All entanglement information is contained in

$$G_0(\vec{x}_s, \vec{y}_s; \vec{x}_i, \vec{y}_i) = e^{-\frac{1}{2}\mathcal{D}_s(|\vec{x}_s - \vec{y}_s|)} e^{-\frac{1}{2}\mathcal{D}_i(|\vec{x}_i - \vec{y}_i|)} \Psi(\vec{x}_s, \vec{x}_i) \Psi^*(\vec{y}_s, \vec{y}_i)$$

Remark:

The turbulent medium induces local decoherence:

- a) pure state \rightarrow mixed state
- b) non-unitary \Rightarrow disentanglement?

Quantification of Entanglement

Biphoton density matrix

$$G_0(\vec{x}_s, \vec{y}_s; \vec{x}_i, \vec{y}_i) = e^{-\frac{1}{2}\mathcal{D}_s(|\vec{x}_s - \vec{y}_s|)} e^{-\frac{1}{2}\mathcal{D}_i(|\vec{x}_i - \vec{y}_i|)} \Psi(\vec{x}_s, \vec{x}_i) \Psi^*(\vec{y}_s, \vec{y}_i)$$

with
$$\Psi(\vec{x}_s, \vec{x}_i) = N \exp\left[-\frac{B}{2}(\vec{x}_s - \vec{x}_i)^2\right] \exp\left[-\frac{A}{2}(\vec{x}_s + \vec{x}_i)^2\right]$$

$$D_s(r) = D_i(r) = \alpha r^{5/3}$$

For CV entanglement

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The second moments of the variables provide useful information about the degree of entanglement. #

Measures of entanglement

- 1. EPR uncertainty
- 2. Entanglement of formation

Generation 2

Biphoton density matrix

 $G_0(\vec{x}_s, \vec{y}_s; \vec{x}_i, \vec{y}_i) = e^{-\frac{1}{2}\mathcal{D}_s(|\vec{x}_s - \vec{y}_s|)} e^{-\frac{1}{2}\mathcal{D}_i(|\vec{x}_i - \vec{y}_i|)} \Psi(\vec{x}_s, \vec{x}_i) \Psi^*(\vec{y}_s, \vec{y}_i)$

with
$$\Psi(\vec{x}_s, \vec{x}_i) = N \exp\left[-\frac{B}{2}(\vec{x}_s - \vec{x}_i)^2\right] \exp\left[-\frac{A}{2}(\vec{x}_s + \vec{x}_i)^2\right]$$
$$D_s(r) = D_i(r) = \alpha r^p, \qquad p = \frac{5}{3}$$

Covariance matrix

$$\begin{split} \Delta^2 x_s &= \Delta^2 x_i = \frac{1}{4} \left(A^{-1} + B^{-1} \right) \\ \Delta^2 p_s &= \Delta^2 p_i = A + B + \alpha p (p-1) \epsilon^{p-2} - \frac{(\alpha p)^2}{2} \epsilon^{2(p-1)}, \quad \epsilon \to 0 \\ \Delta x_s \Delta x_i &= \frac{1}{4} \left(A^{-1} - B^{-1} \right) \\ \Delta p_s \Delta p_i &= A - B \\ \Delta x_s \Delta p_s &= \Delta x_i \Delta p_i = \Delta x_s \Delta p_i = \Delta x_i \Delta p_s = 0 \end{split}$$

Generation Gaussian Approximation 2

Biphoton density matrix

 $G_0(\vec{x}_s, \vec{y}_s; \vec{x}_i, \vec{y}_i) = e^{-\frac{1}{2}\mathcal{D}_s(|\vec{x}_s - \vec{y}_s|)} e^{-\frac{1}{2}\mathcal{D}_i(|\vec{x}_i - \vec{y}_i|)} \Psi(\vec{x}_s, \vec{x}_i) \Psi^*(\vec{y}_s, \vec{y}_i)$

with
$$\Psi(\vec{x}_s, \vec{x}_i) = N \exp\left[-\frac{B}{2}(\vec{x}_s - \vec{x}_i)^2\right] \exp\left[-\frac{A}{2}(\vec{x}_s + \vec{x}_i)^2\right]$$
$$D_s(r) = D_i(r) = \alpha r^p, \qquad p = \frac{5}{3}$$

Covariance matrix

$$\Delta^{2} x_{s} = \Delta^{2} x_{i} = \frac{1}{4} \left(A^{-1} + B^{-1} \right)$$

$$\Delta^{2} p_{s} = \Delta^{2} p_{i} = A + B + \alpha p (p-1) \epsilon^{p-2} - \frac{(\alpha p)^{2}}{2} \epsilon^{2(p-1)}, \quad \epsilon \to 0$$

$$\Delta x_{s} \Delta x_{i} = \frac{1}{4} \left(A^{-1} - B^{-1} \right)$$

$$\Delta p_{s} \Delta p_{i} = A - B$$
Singularity in $\Delta^{2} p_{s}$

$$\Delta x_{s} \Delta p_{s} = \Delta x_{i} \Delta p_{i} = \Delta x_{s} \Delta p_{i} = \Delta x_{i} \Delta p_{s} = 0$$

Gaussian Approximation 2

Biphoton density matrix

$$G_{0}(\vec{x}_{s}, \vec{y}_{s}; \vec{x}_{i}, \vec{y}_{i}) = e^{-\frac{1}{2}\mathcal{D}_{s}(|\vec{x}_{s} - \vec{y}_{s}|)} e^{-\frac{1}{2}\mathcal{D}_{i}(|\vec{x}_{i} - \vec{y}_{i}|)} \Psi(\vec{x}_{s}, \vec{x}_{i}) \Psi^{*}(\vec{y}_{s}, \vec{y}_{i})$$
with $\Psi(\vec{x}_{s}, \vec{x}_{i}) = N \exp\left[-\frac{B}{2}(\vec{x}_{s} - \vec{x}_{i})^{2}\right] \exp\left[-\frac{A}{2}(\vec{x}_{s} + \vec{x}_{i})^{2}\right]$

$$D_s(r) = \alpha r^{5/3} \longrightarrow \alpha r^{6/3}$$



Expect their impacts to entanglement degradation are very similar.

Gaussian Approximation 2

Biphoton density matrix

$$G_0(\vec{x}_s, \vec{y}_s; \vec{x}_i, \vec{y}_i) = e^{-\frac{1}{2}\mathcal{D}_s(|\vec{x}_s - \vec{y}_s|)} e^{-\frac{1}{2}\mathcal{D}_i(|\vec{x}_i - \vec{y}_i|)} \Psi(\vec{x}_s, \vec{x}_i) \Psi^*(\vec{y}_s, \vec{y}_i)$$

with
$$\Psi(\vec{x}_{s}, \vec{x}_{i}) = N \exp\left[-\frac{B}{2}(\vec{x}_{s} - \vec{x}_{i})^{2}\right] \exp\left[-\frac{A}{2}(\vec{x}_{s} + \vec{x}_{i})^{2}\right]$$

$$D_s(r) = \alpha r^{5/3} - OK \quad \alpha r^{6/3}$$

Covariance matrix

$$\Delta^2 x_s = \Delta^2 x_i = \frac{1}{4} \left(A^{-1} + B^{-1} \right)$$
$$\Delta^2 p_s = \Delta^2 p_i = A + B + 2\alpha$$
$$\Delta x_s \Delta x_i = \frac{1}{4} \left(A^{-1} - B^{-1} \right)$$
$$\Delta p_s \Delta p_i = A - B$$
$$\Delta x_s \Delta p_s = \Delta x_i \Delta p_i = \Delta x_s \Delta p_i = \Delta x_i \Delta p_s = 0$$

No singularity

Phase structure function (Kolmogorov + Gaussian approx.)

Control Parameters

$$D_s(r) = 3.44 \left(\frac{r}{r_0}\right)^{6/3}$$

 $D = 2\Delta x_s = 2\Delta x_i = \sqrt{A^{-1} + B^{-1}}$

 r_0 – the length scale of turbulent structure

D

Effective aperture diameter:

The Institute of W

Control p

parameters – (i) degree of turbulence:
$$\frac{D}{r_0}$$

(ii) degree of initial entanglement: $\eta = \frac{B}{A}$
 $\Psi(\vec{x}_s, \vec{x}_i) = N \exp\left[-\frac{B}{2}(\vec{x}_s - \vec{x}_i)^2\right] \exp\left[-\frac{A}{2}(\vec{x}_s + \vec{x}_i)^2\right]$

Effect of Turbulence to Entanglement

EPR uncertainty

$$\Delta = \sqrt{\Delta^2 (x_s - x_i) + \Delta^2 (p_s + p_i)}$$

In general: $\Delta < 1$ entangled

$$\Delta \geq 1$$
 entangled / disentangled

Giedke et al., PRL 91, 107901 (2003)

Effect of Turbulence to Entanglement

EPR uncertainty

$$\Delta = \sqrt{\Delta^2 (x_s - x_i) + \Delta^2 (p_s + p_i)}$$

Gaussian state:

 $\Delta < 1$ entangled $\Delta > 1$ disentangled

We find #
$$\Delta = \sqrt{\frac{(1+\eta^{-1}) + 3.44 (D/r_0)^2}{1+\eta}}$$



Note:

1. Entanglement of the two photons is totally destroyed if $3.44 \left(\frac{D}{r_0}\right)^2 \geq \eta - \eta^{-1}$

2.
$$2\log \Delta = \log \left[1 + \frac{3.44}{1+\eta^{-1}} \left(\frac{D}{r_0}\right)^2 \right] + \log \left(\frac{1+\eta^{-1}}{1+\eta}\right)$$

Giedke et al., PRL 91, 107901 (2003)

Effect of Turbulence to Entanglement

EPR uncertainty

 $\Delta <$

$$\Delta = \sqrt{\Delta^2 (x_s - x_i) + \Delta^2 (p_s + p_i)}$$

Gaussian state:

$$\Delta < 1$$
 entangled $\Delta \ge 1$ disentangled

We find #
$$\Delta = \sqrt{\frac{(1+\eta^{-1}) + 3.44 (D/r_0)^2}{1+\eta}}$$

Entanglement of formation for Gaussian states

how much entanglement is needed to _ construct the state

$$E_F = c_+ \log c_+ - c_- \log c_-$$

where $c_\pm = rac{1}{4} \left(\Delta^{-1/2} \pm \Delta^{1/2}
ight)^2$

Giedke et al., PRL 91, 107901 (2003)



Experiment & Preliminary Results

Turbulence medium:

heat gun (easy to implement)



The Institute of **Experiment & Preliminary Results**



Summary / Outlook

<u>Summary</u>

- Constructed the theory of SPDC spatial coordinates with turbulence effect
- Approximated the system by a Gaussian model:
 - *initial wave function* (position & momentum coordinates)
 - *turbulence power law* (phase structure function: $5/3 \rightarrow 6/3$)
- Shown effect of turbulence of the entanglement. Entanglement is strongly affected at $D/r_0 > 0.5$.
- Obtained preliminary experimental results

Next steps

- Quantitative measurement of entanglement vs. turbulence effect
- Use adaptive optics to minimize disentanglement





Thank you

Q & A

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