## "Slow" Light in Nanostructured Devices Robert W. Boyd with John Heebner, Nick Lepeshkin, Aaron Schweinsberg, and Q-Han Park

The Institute of Optics, University of Rochester, Rochester, NY 14627

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## NLO of SCISSOR Devices

(Side-Coupled Integrated Spaced Sequence of Resonators)



Shows slow-light, tailored dispersion, and enhanced nonlinearity Optical solitons described by nonlinear Schrodinger equation

• Weak pulses spread because of dispersion



• But intense pulses form solitons through balance of dispersion and nonlinearity.



Ultrafast All-Optical Switch Based On Arsenic Triselenide Chalcogenide Glass

• We excite a whispering gallery mode of a chalcogenide glass disk.



- The nonlinear phase shift scales as the square of the finesse F of the resonator. (F  $\approx 10^2$  in our design)
- Goal is 1 pJ switching energy at 1 Tb/sec.



J. E. Heebner and R. W. Boyd, Opt. Lett. 24, 847, 1999. (implementation with Dick Slusher, Lucent)

## **A Real Whispering Gallery**



## St. Paul's Cathedral, London

#### **Alliance for Nanomedical Technologies**

## **Photonic Devices for Biosensing**

## **Objective:**

Obtain high sensitivity, high specificity detection of pathogens through optical resonance

## Approach:

Utilize high-finesse whispering-gallerymode disk resonator.

Presence of pathogen on surface leads to dramatic decrease in finesse.

## Simulation of device operation:



Intensity distribution in absense of absorber.



Intensity distribution in presence of absorber.



Slow Light, Induced Dispersion, Enhanced Nonlinearity, Self-Steepening, and Optical Solitons in a Side-Coupled Integrated Spaced Sequence Of Resonators



# SCISSOR

John E. Heebner Q-Han Park Robert W. Boyd



# Motivation

To exploit the ability of microresonators to enhance nonlinearities and induce strong dispersive effects for creating structured waveguides with exotic properties.

Currently, most of the work done in microresonators involves applications such as disk lasers, dispersion compensators and add-drop filters. There's not much nonlinear action!

A cascade of resonators side-coupled to an ordinary waveguide can exhibit:



- slow light propagation
- induced dispersion
- enhanced nonlinearities



# **Properties of a Single Microresonator**

Assuming negligible attenuation, this resonator is, unlike a Fabry-Perot, of the "all-pass" device there is no reflected or drop port.



 $E_2$ 

 $E_1$ 



Build-up Factor Intensity Enhancement ( $|E_3 / E_1|^2$ )  $(E_1)^2$  $r^2 = 0.90$ **Definitions**  $r^2 = 0.75$ Щ  $r^2 = 0.25$  $r^2 = 0.00$ **Finesse**  $F = \frac{\pi}{1-r}$ Modified Dispersion Relation ( $\beta$  vs.  $\omega$ ) effective propagation  $r^2 = 0.00$ constant ( $\beta$ ) Transit Time  $r^2 = 0.75$  $\underline{n2\pi R}$  $r^2 = 0.25$  $r^2 = 0.90$  $\omega_{\rm R} + \frac{2\pi}{T}$  $\omega_{\rm R}$  $\omega_{\rm R} - \frac{2\pi}{T}$ frequency (m)

# **Propagation Equation for a SCISSOR**



By arranging a spaced sequence of resonators, side-coupled to an ordinary waveguide, one can create an effective, structured waveguide that supports pulse propagation in the NLSE regime.

Propagation is unidirectional, and there is NO photonic bandgap to produce the enhancement. Feedback is intra-resonator and not inter-resonator.

> Nonlinear Schrödinger Equation (NLSE)  $\frac{\partial}{\partial z}A = -i\frac{1}{2}\beta_2\frac{\partial^2}{\partial t^2}A + i\gamma|A|^2A$ Fundamental Soliton Solution  $A(z,t) = A_0 \operatorname{sech}\left(\frac{t}{T_p}\right)e^{i\frac{1}{2}\gamma|A_0|^2z}$

# **Balancing Dispersion & Nonlinearity**





Resonator-induced dispersion can be 5-7 orders of magnitude greater than the material dispersion of silica!

Resonator enhancement of nonlinearity can be 3-4 orders of magnitude!

An enhanced nonlinearity may be balanced by an induced anomalous dispersion at some detuning from resonance to form solitons

A characteristic length, the soliton period may as small as the distance between resonator units!

# **Soliton Propagation**



## **Dark Solitons**

SCISSOR system also supports the propagation of dark solitons.



## **Higher-Order Effects - Self-Steepening**

Higher order dispersive terms such as  $\beta_3$  are present in the system and become more dominant as the pulsewidth becomes nearly as short as the cavity lifetime. Because the nonlinear enhancement is in fact frequency dependent, or (equivalently here) because the group velocity is intensity dependent, selfsteepening of pulses is possible even for relatively long pulse widths.

#### A generalized NLSE:

$$\frac{\partial}{\partial z}A = -i\frac{1}{2}\beta_2\frac{\partial^2}{\partial t^2}A - \frac{1}{6}\beta_3\frac{\partial^3}{\partial t^3}A + i\gamma|A|^2A - s\frac{\partial}{\partial t}|A|^2A$$

Self-steepening of a 20 ps Gaussian pulse after 100 resonators



# **Soliton Splitting and Compression**

The dispersive nature of the nonlinear enhancement (self-steepening) leads to an intensity-dependent group velocity which splits an N-order soliton into N fundamental solitons of differing peak intensities and widths.

Here, a 2<sup>nd</sup> - order "breathing" soliton splits into 2 fundamental solitons:



# Limitations

#### Bandwidth

A 5  $\mu$ m diamater disk with a finesse of 200 possesses a bandwidth of 100 GHz Cascading N resonators further decreases the available bandwidth; fortunately the scaling is governed by  $\beta_4$  effects which results in a reduction of N<sup>1/4</sup>



#### Manufacturability

Tolerances on resonator diameter and gap spacing to be met within 200 nm variance in gap spacing analogous to hyperfine distribution variance in radius/index analogous to Doppler broadening

Attenuation - dominated by scattering losses may be counterbalanced by:

1) gain

2) dispersion decreasing SCISSOR :



## **Radiation Losses**

Is it really possible to confine light within a wavelength-size structure?

- Radiation losses provide fundamental limit to the finesse of a WGM resonator.
- Radiation losses are analogous to fiber bending losses.
- But these losses can be rendered negligible through use of a large dielectric contrast.



## Conclusions

The SCISSOR exhibits strong dispersive & nonlinear properties and supports solitons. It incorporates the essence of photonic bandgap (PBG) enhancement found in structures such as CROWs & Bragg gratings *without the gap itself*.

Manufacturability will require some ingenuity, but already resonators are already in use as all-pass filters and tunable dispersion

SCISSORs have the potential for producing soliton propagation effects at the integrated photonics scale. The possibilities for structurally engineered waveguides of this type include:

- Pulse compression / imaging in an integrated device
- Optical Time Division Multiplexing (OTDM)
- Soliton-based optical switching (perhaps w/ a few photons)
- Slow-light propagation (group-matched acousto-optics)
- EO,TO Tunable NLSE propagation (A theorist's dream)
- Just about any other generalized NLSE effect

## Nanofabrication

- Materials (artificial materials)
- Devices

(distinction?)

#### Microdisk Resonator Design

(Not drawn to scale) All dimensions in microns



### **Photonic Device Fabrication Procedure**



RWB - 10/4/01

## **Nonlinear Optical Loop-De-Loop**



J.E. Heebner and R.W.B.

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## Photonic Devices in GaAs/AlGaAs





## **Microdisk Resonator Design / Vertically Coupled Structure**

