

# Slow, Fast, and “Backwards” Light: Fundamental Aspects

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# Interest in Slow Light

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Intrigue: Can (group) refractive index really be  $10^6$ ?

Fundamentals of optical physics

Optical delay lines, optical storage, optical memories

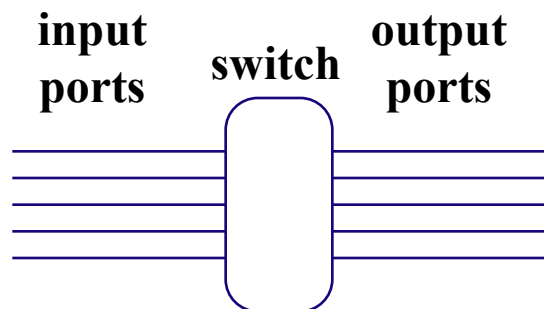
Implications for quantum information

What about fast light ( $v > c$ ) and backwards light ( $v$  negative)?

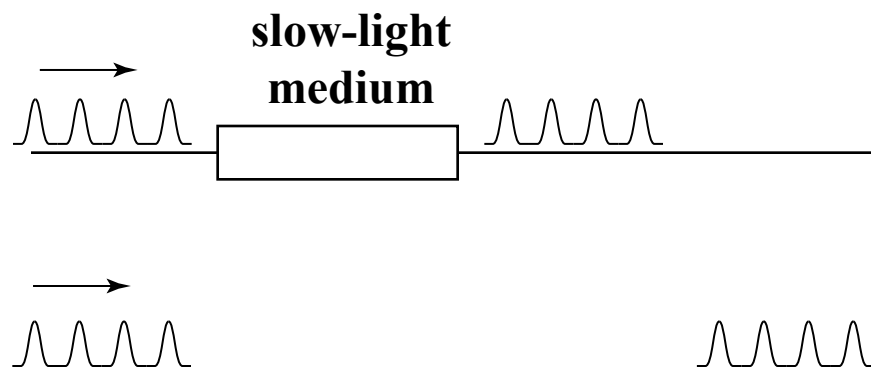
Boyd and Gauthier, "Slow and Fast Light," in Progress in Optics, 43, 2002.



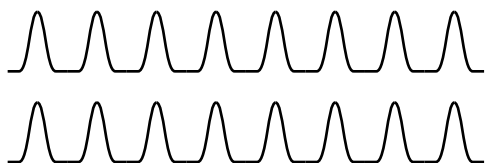
## All-Optical Switch



## Use Optical Buffering to Resolve Data-Packet Contention



**But what happens if two data packets arrive simultaneously?**



**Controllable slow light for optical buffering can dramatically increase system performance.**

# Slow and Fast Light and Optical Resonances

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Pulses propagate at the group velocity given by

$$v_g = \frac{c}{n_g} \quad n_g = n + \omega \frac{dn}{d\omega}$$

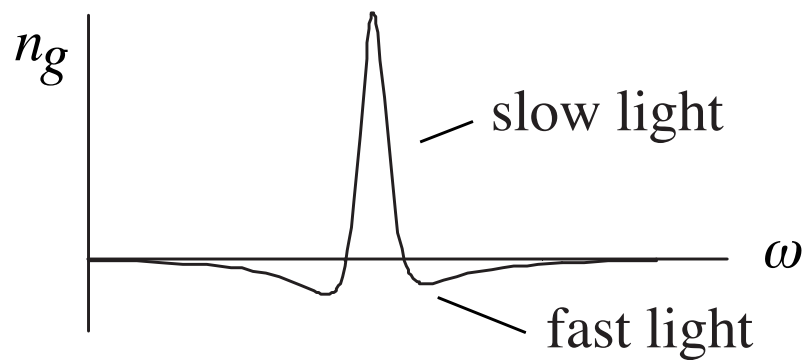
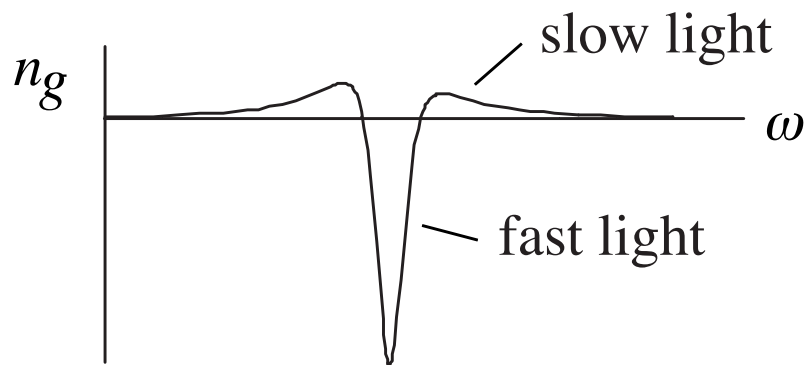
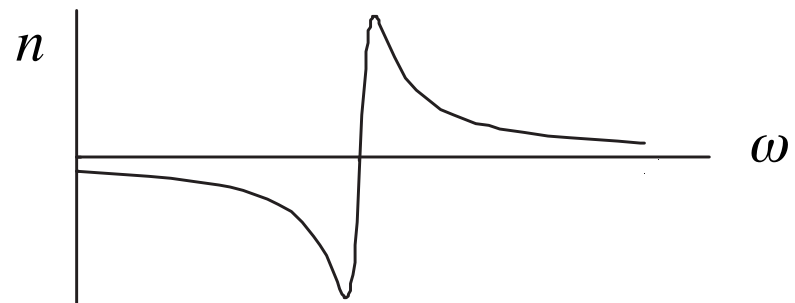
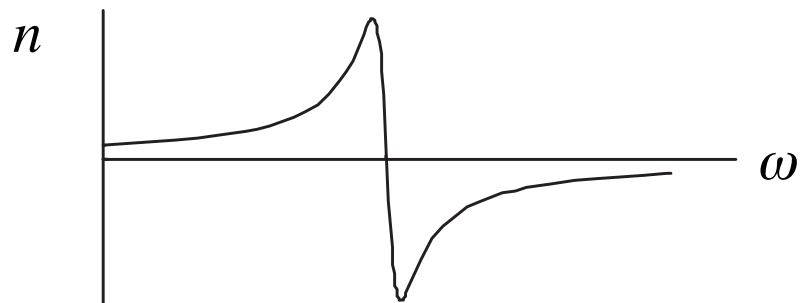
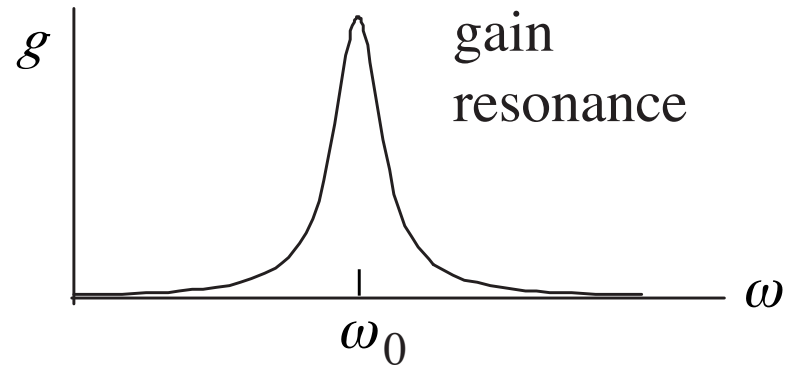
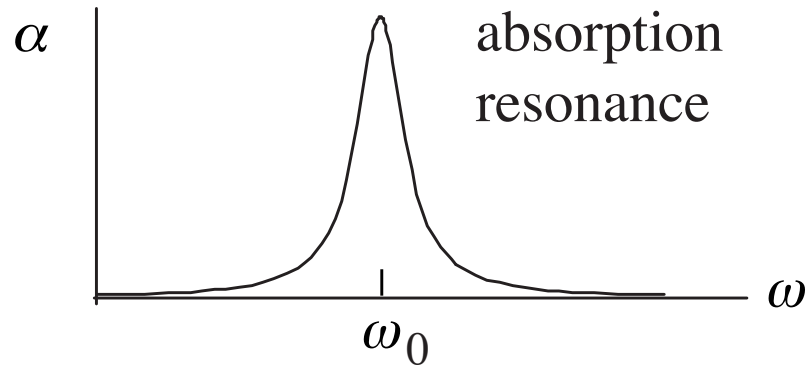
Want large dispersion to obtain extreme group velocities

Sharp spectral features produce large dispersion.

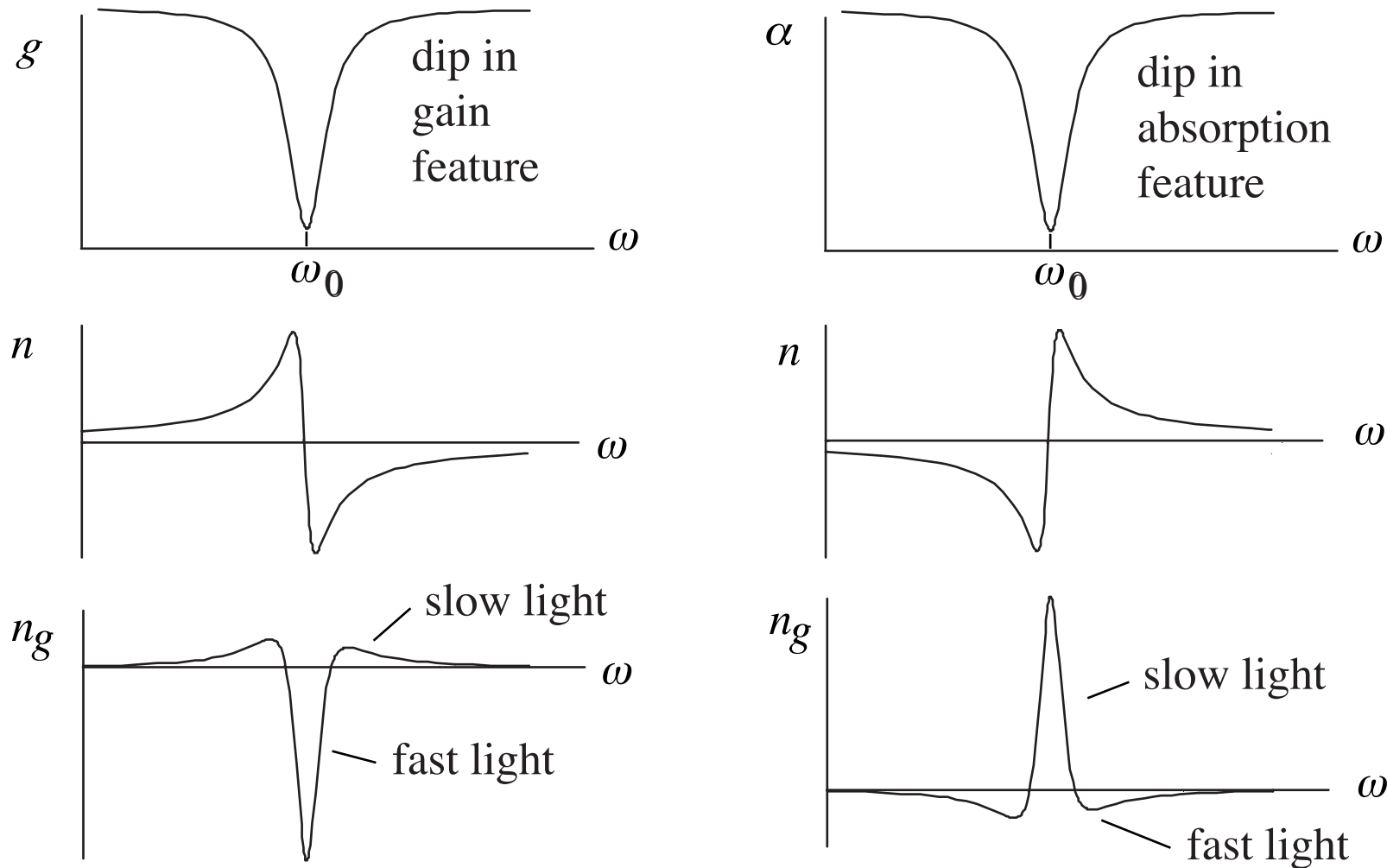
The group index can be large and positive (slow light), positive and much less than unity (fast light) or negative (backwards light).

# How to Create Slow and Fast Light I – Use Isolated Gain or Absorption Resonance

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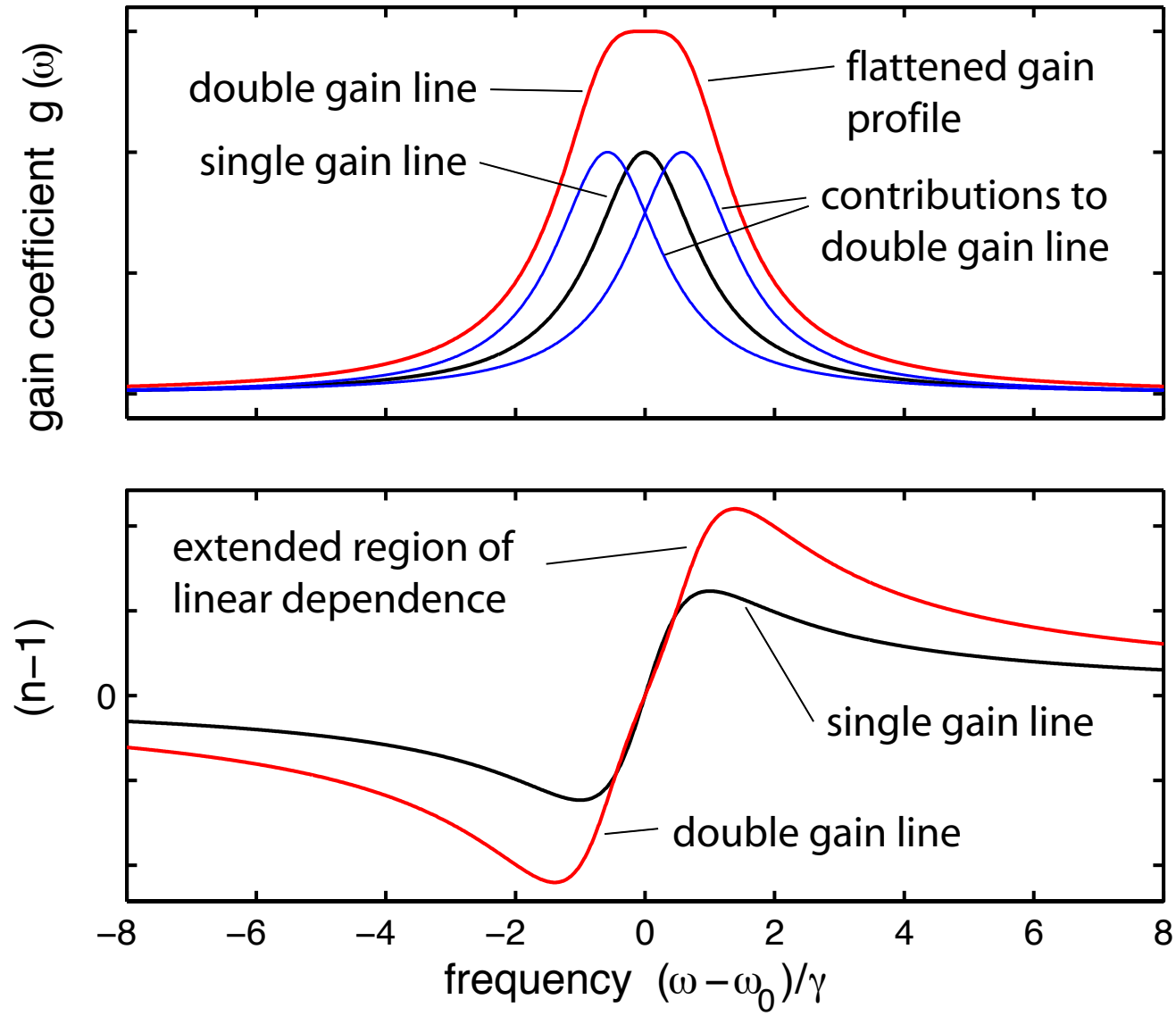


# How to Create Slow and Fast Light II – Use Dip in Gain or Absorption Feature



Narrow dips in gain and absorption lines can be created by various nonlinear optical effects, such as electromagnetically induced transparency (EIT), coherent population oscillations (CPO), and conventional saturation.

# How to Create Slow and Fast Light III – Dispersion Management



# Numerical Modeling of Pulse Propagation through Slow and Fast-Light Media

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Numerically integrate the reduced wave equation

$$\frac{\partial A}{\partial z} - \frac{1}{v_g} \frac{\partial A}{\partial t} = 0$$

and plot  $A(z,t)$  versus distance  $z$ .

Assume an input pulse with a Gaussian temporal profile.

Study three cases:

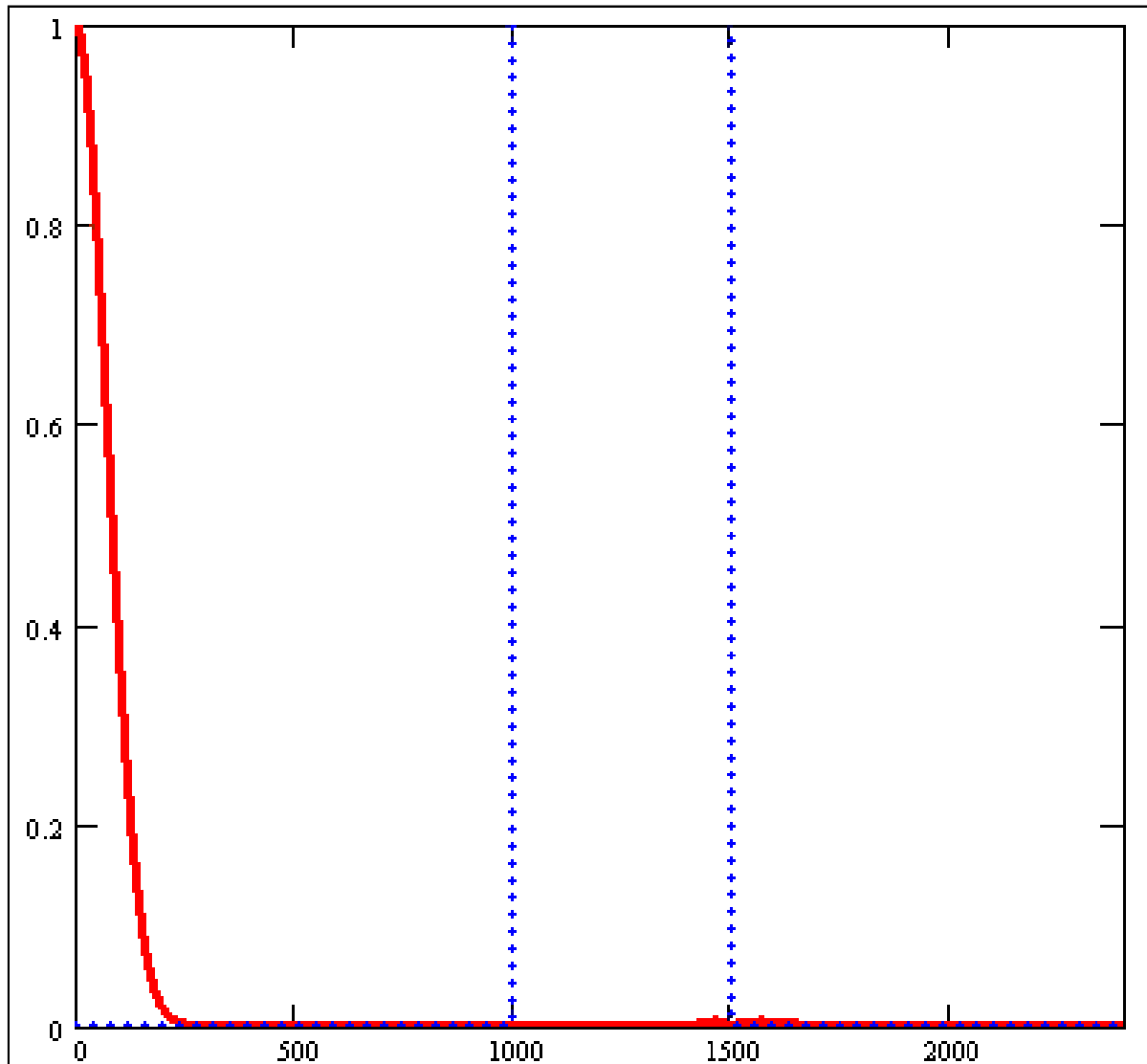
Slow light  $v_g = 0.5 c$

Fast light  $v_g = 5 c$  and  $v_g = -2 c$

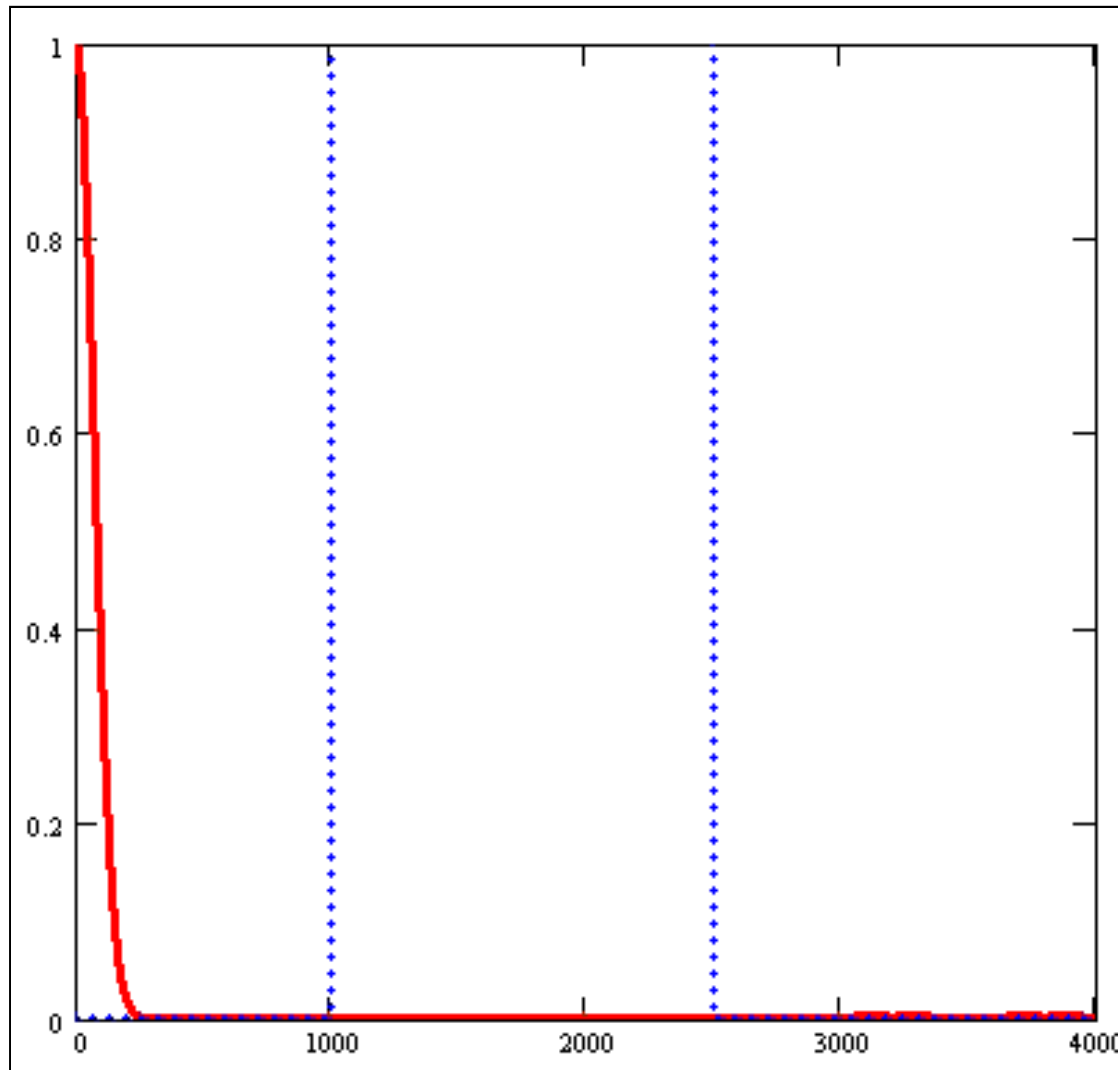
**CAUTION:** This is a very simplistic model. It ignores GVD and spectral reshaping.



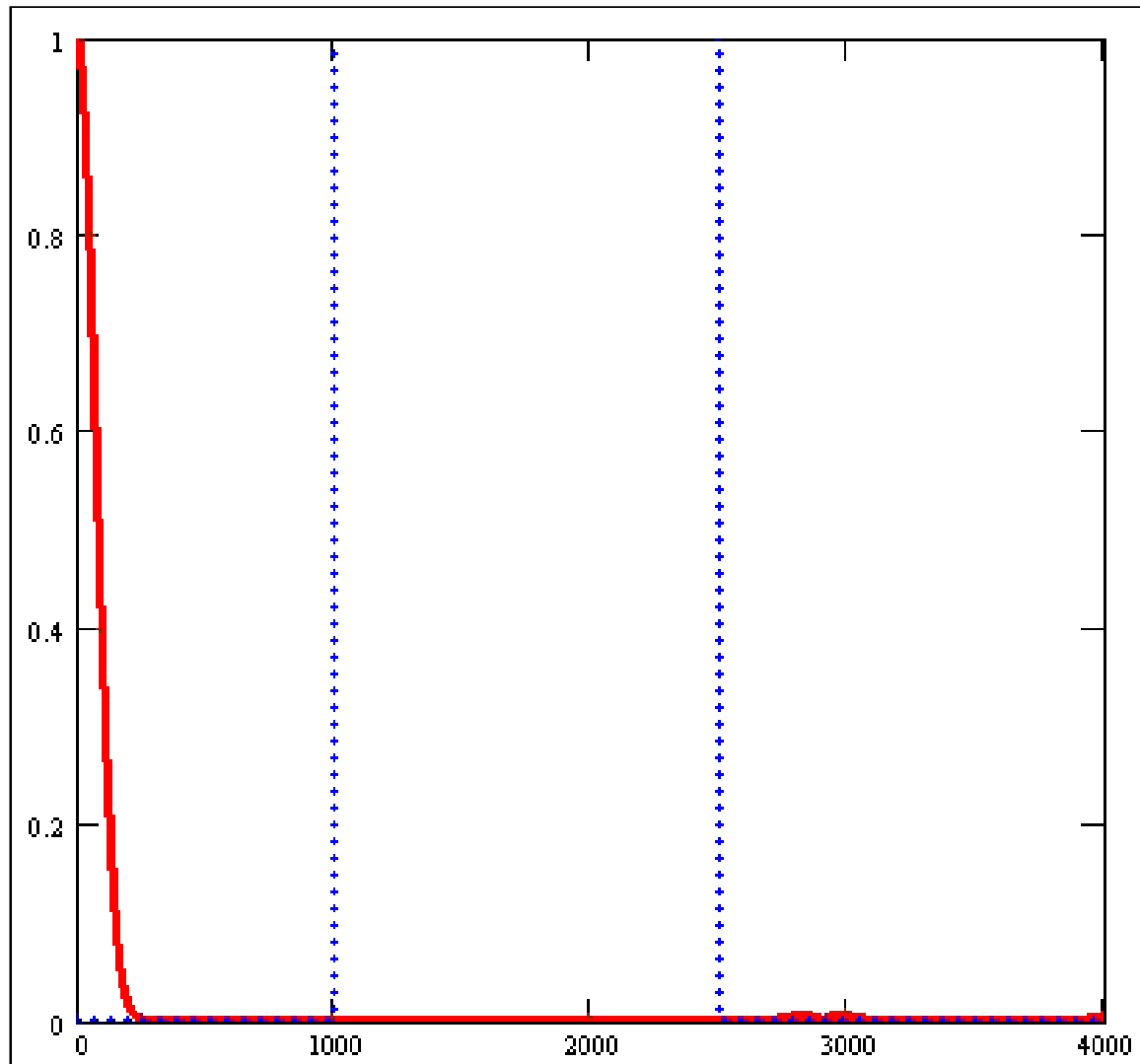
# Pulse Propagation through a Slow-Light Medium ( $n_g = 2$ , $v_g = 0.5 c$ )



# Pulse Propagation through a Fast-Light Medium ( $n_g = .2$ , $v_g = 5 c$ )

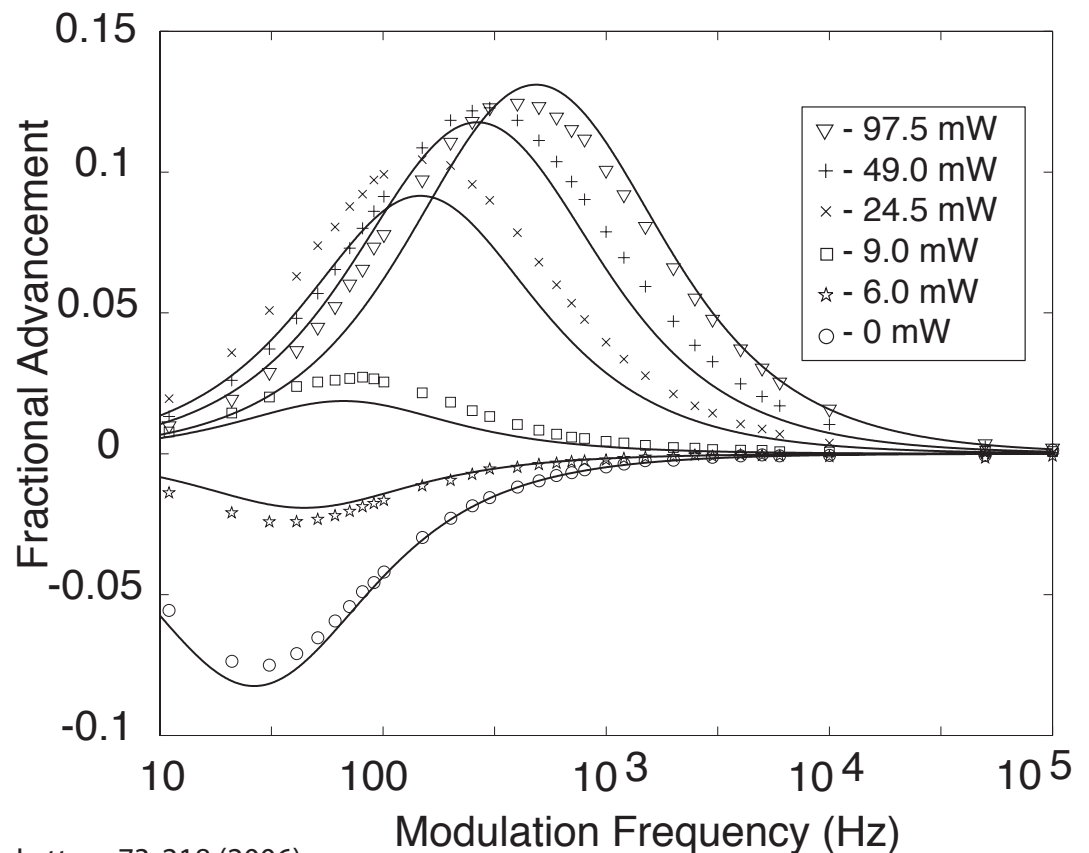
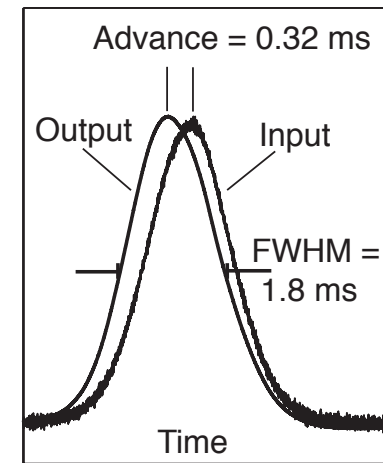
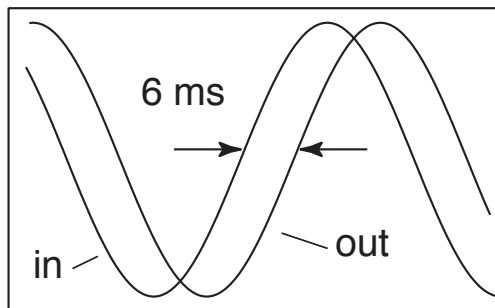
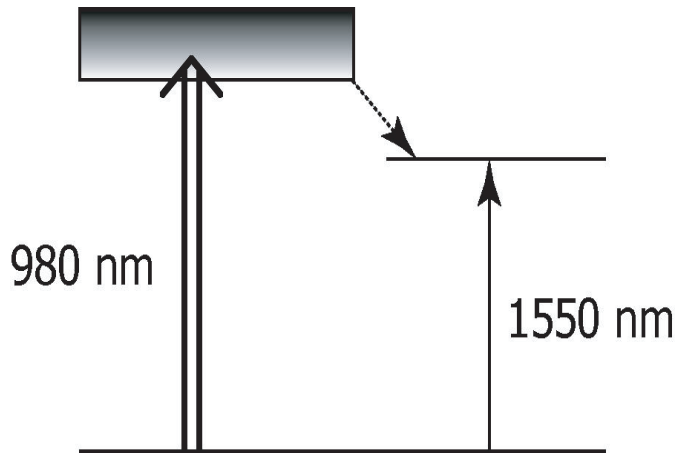


# Pulse Propagation through a Fast-Light Medium ( $n_g = -.5$ , $v_g = -2 c$ )

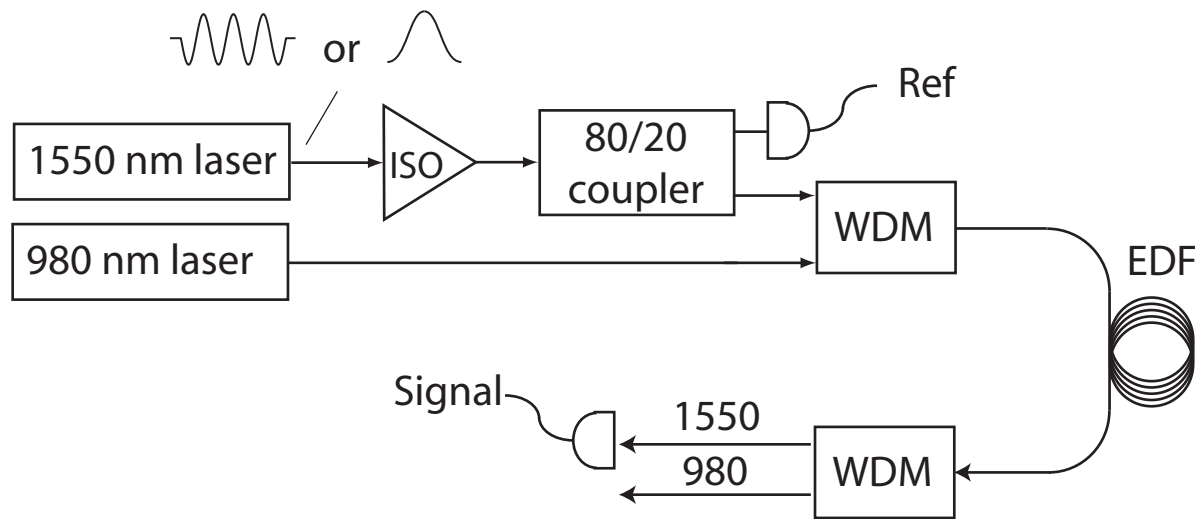


# Slow and Fast Light in an Erbium Doped Fiber Amplifier

- Fiber geometry allows long propagation length
- Saturable gain or loss possible depending on pump intensity



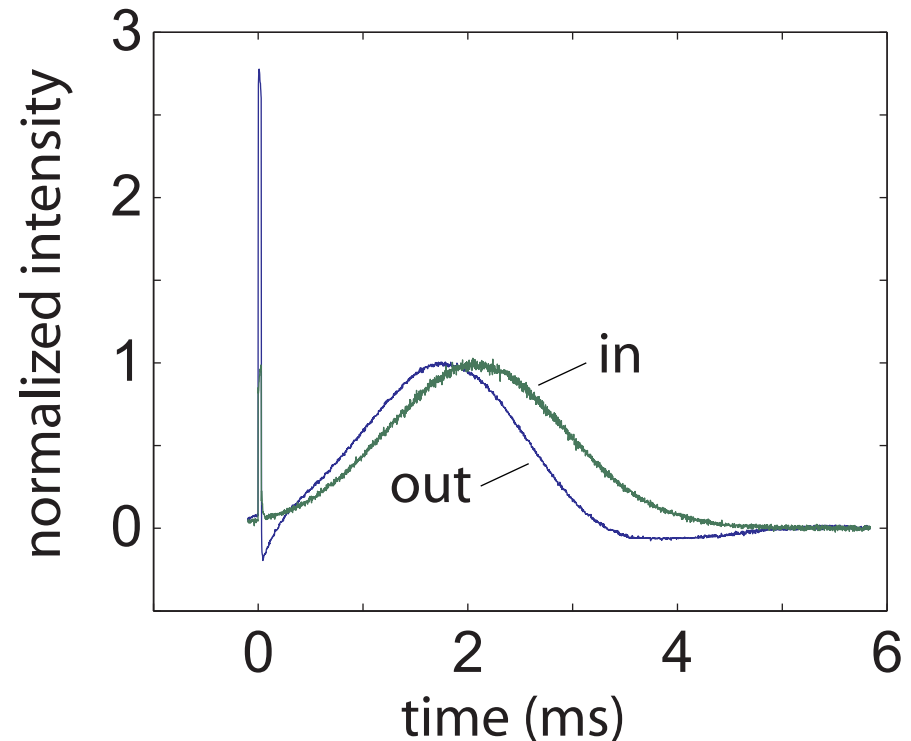
# Observation of Backward Pulse Propagation in an Erbium-Doped-Fiber Optical Amplifier



We time-resolve the propagation of the pulse as a function of position along the erbium-doped fiber.

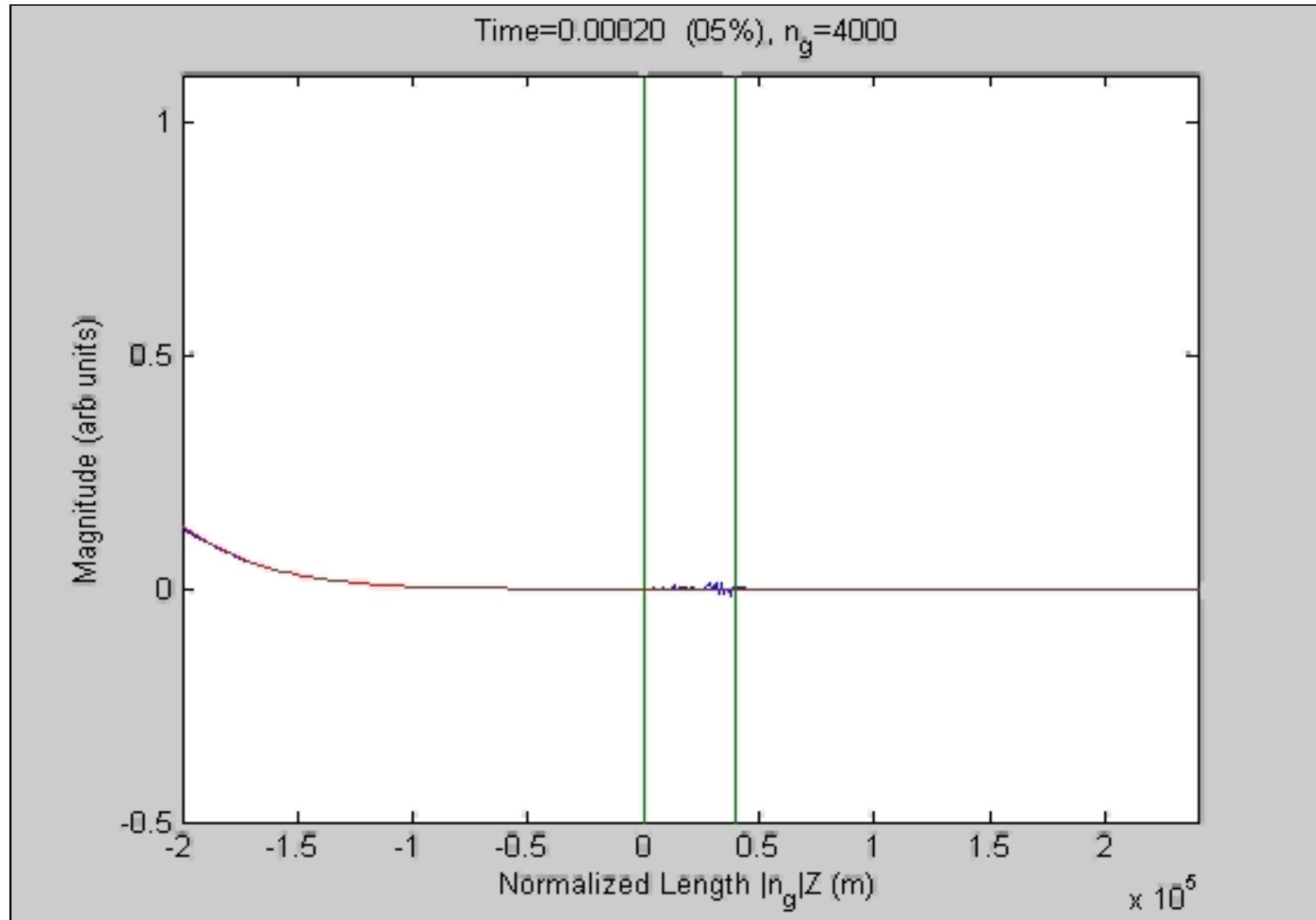
## Procedure

- cutback method
- couplers embedded in fiber



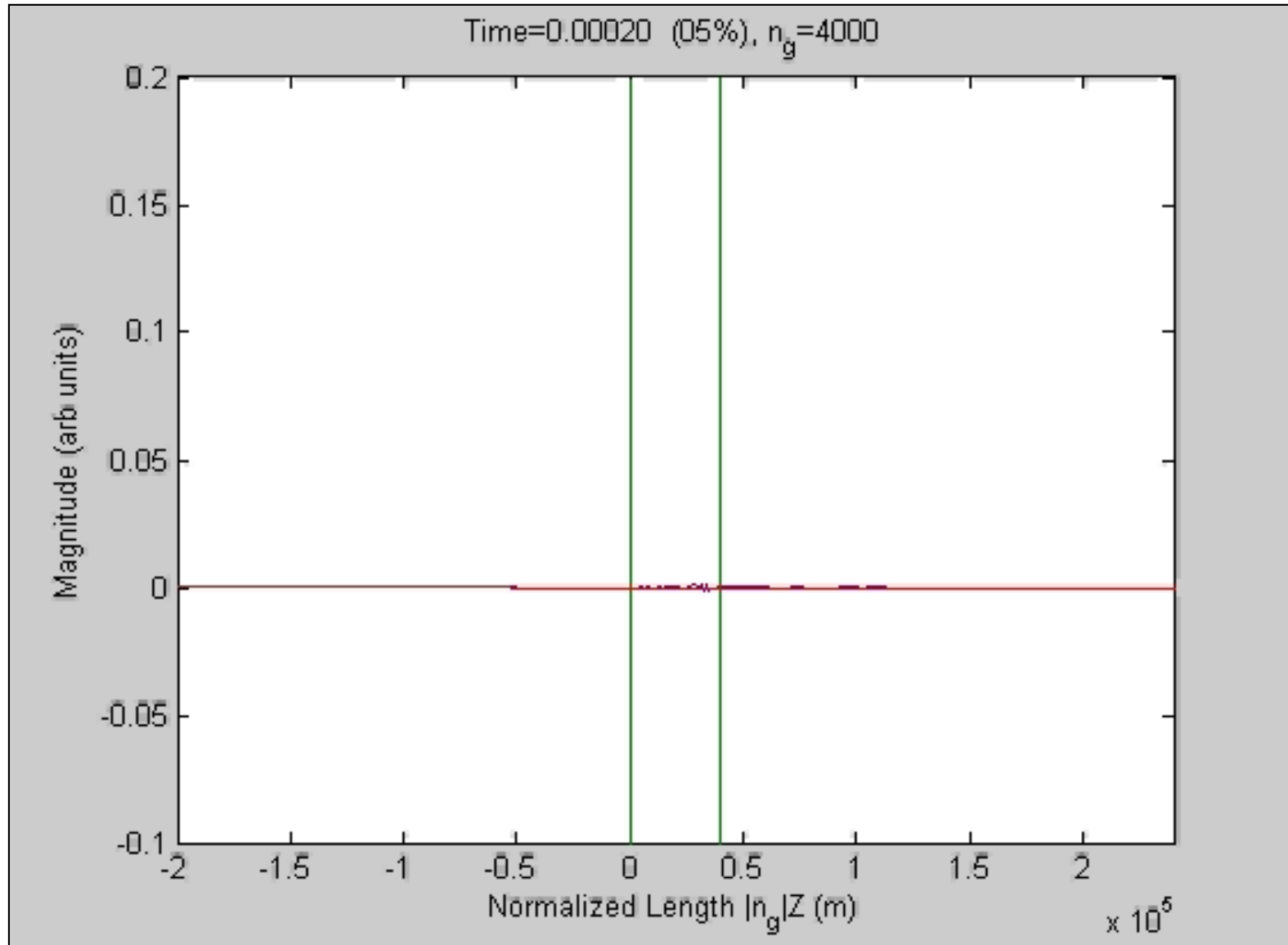
# Experimental Results: Backward Propagation in Erbium-Doped Fiber

Normalized: (Amplification removed numerically)

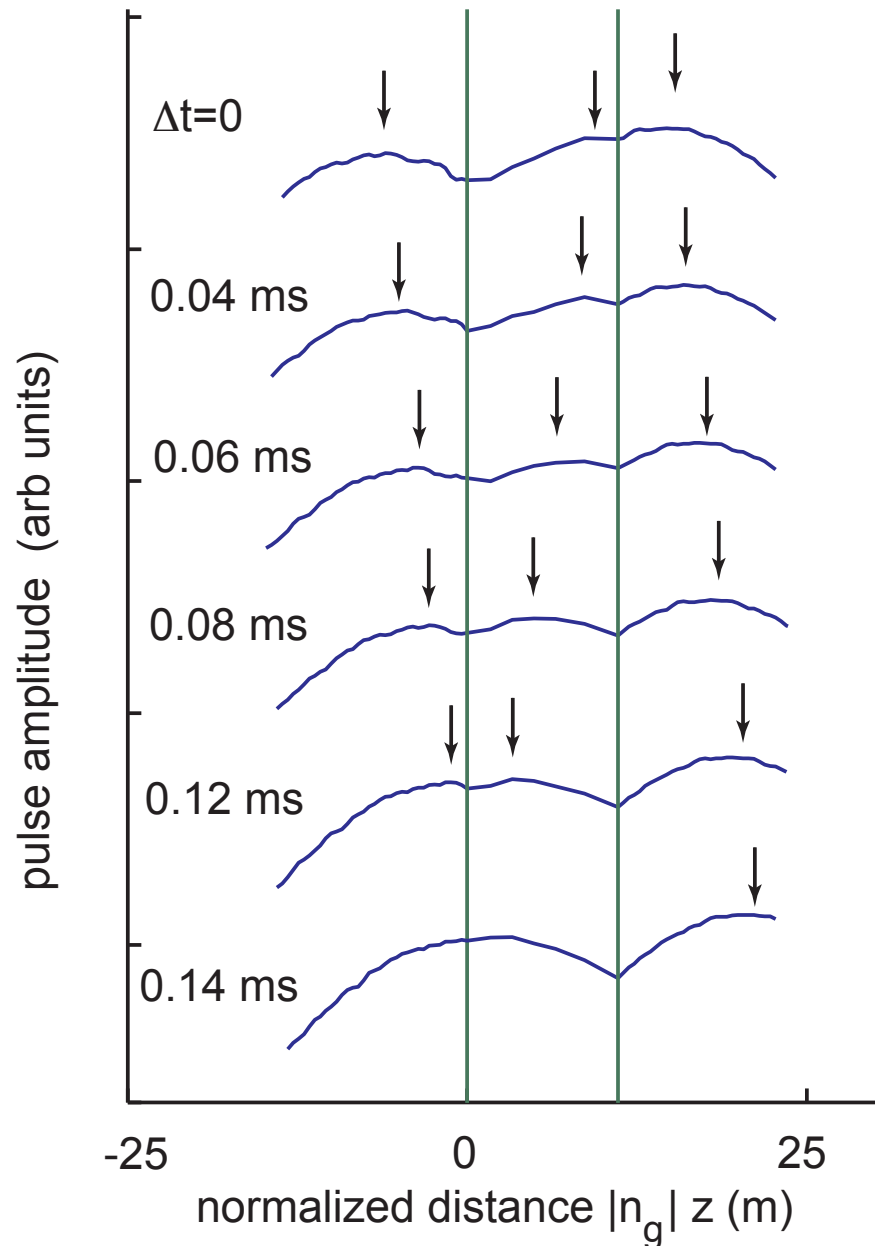


# Experimental Results: Backward Propagation in Erbium-Doped Fiber

## Un-Normalized



# Observation of Backward Pulse Propagation in an Erbium-Doped-Fiber Optical Amplifier





# Observation of Backward Pulse Propagation in an Erbium-Doped-Fiber Optical Amplifier

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Summary:

“Backwards” propagation is a realizable physical effect.

(Of course, many other workers have measured negative time delays. Our contribution was to measure the pulse evolution within the material medium.)

# Fundamental Limits on Slow and Fast Light

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Slow Light: There appear to be no fundamental limits on how much one can delay a pulse of light (although there are very serious practical problems).\*

Fast Light: But there do seem to be essentially fundamental limits to how much one can advance a pulse of light.

**Why are the two cases so different?\***

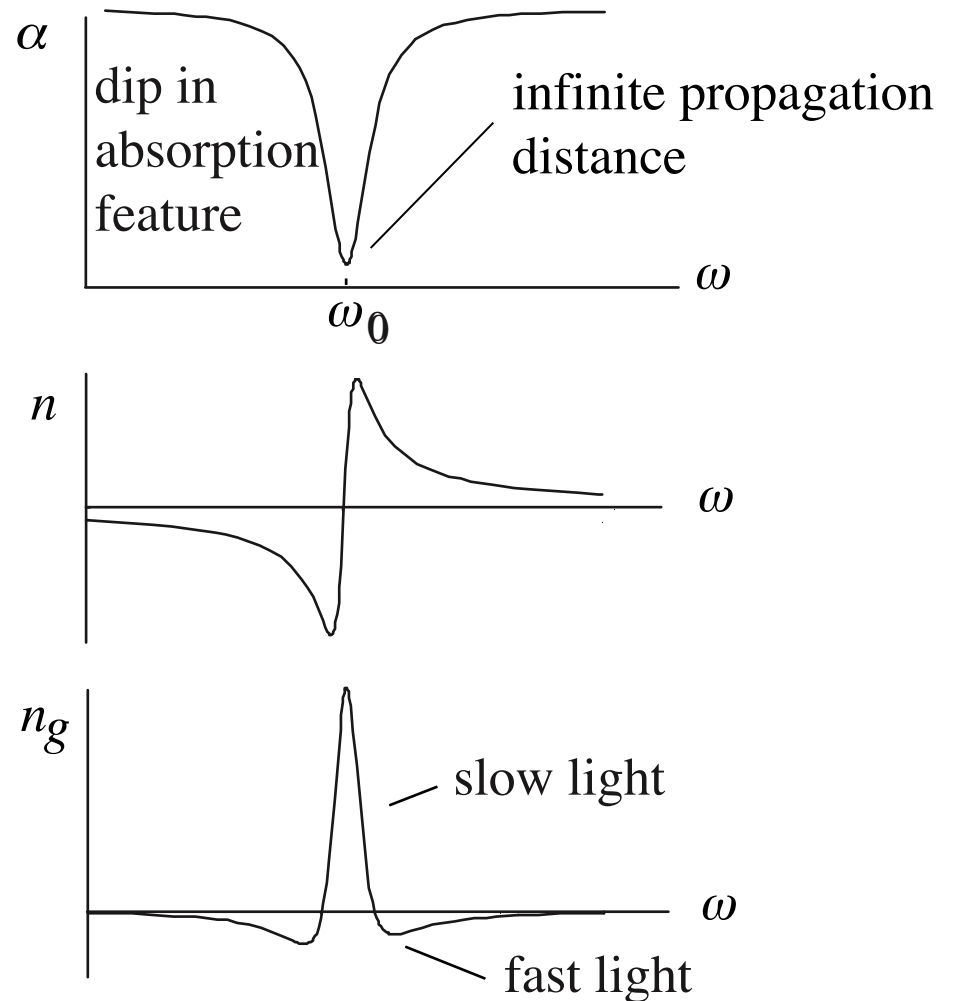
\* Boyd, Gauthier, Gaeta, and Willner, PRA 2005

\*\* We cannot get around this problem simply by invoking causality, first because we are dealing with group velocity (not information velocity), and second because the relevant equations superficially appear to be symmetric between the slow- and fast-light cases.

# Why is there no limit to the amount of pulse delay?

At the bottom of the dip in the absorption, the absorption can in principle be made to vanish. There is then no limit on how long a propagation distance can be used.

This “trick” works only for slow light.



# Why can one delay (but not advance) a pulse by an arbitrarily large amount?

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Two crucial differences between slow and fast light

(1) First, note that we cannot use gains greater than approximately  $\exp(16)$  at any frequency to avoid ASE. And we cannot have absorption larger than  $T = \exp(-16)$  at the signal frequency, so signal can be measured. (Of course, the argument does not hinge on the value 16.) When examined quantitatively, these constraints impose a limit of at most two pulse-widths of delay or advancement.

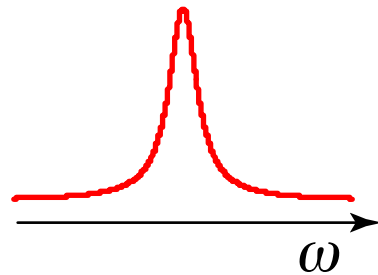
One can overcome these constraints by using a deep hole in an absorption feature, but this trick works only for slow light, as we have just seen.

(2) Spectral reshaping of the pulse is the dominant competing effect in most slow/fast light systems. This also behaves differently for slow and fast-light systems, as we shall now see.

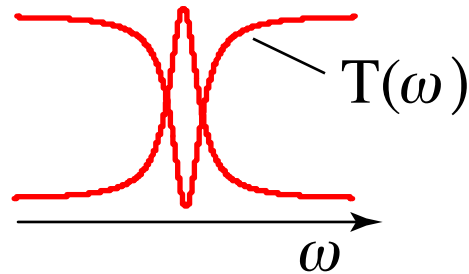
# Influence of Spectral Reshaping (Line-Center Operation, Dip in Gain or Absorption Feature)

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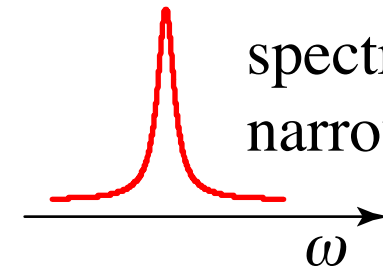
input pulse



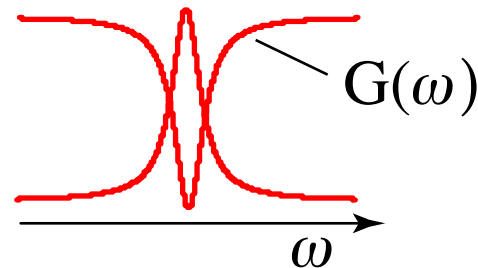
output pulse  
slow-light



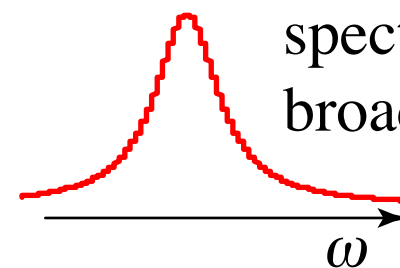
spectrally  
narrowed pulse



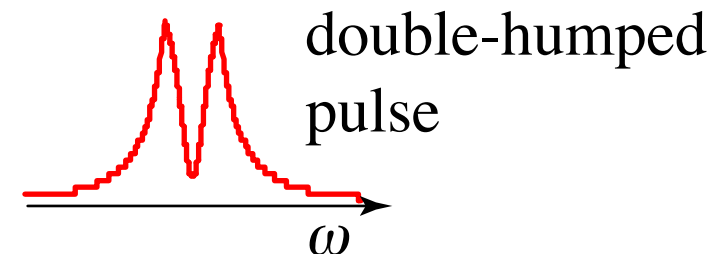
output pulse  
fast-light



spectrally  
broadened pulse



for still longer propagation  
distances, the pulse breaks  
up spectrally and temporally



# Numerical Results: Propagation through a Linear Dispersive Medium

Fast light: single Lorentzian absorption line

$$T = \exp(-16)$$

vary absorption line width to control advance

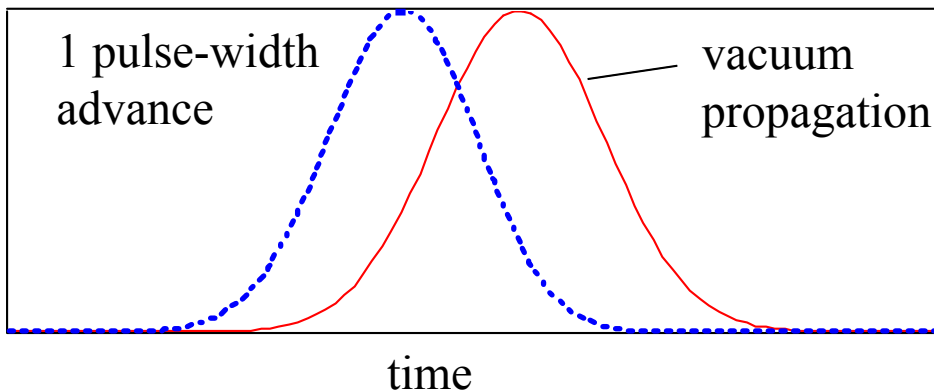
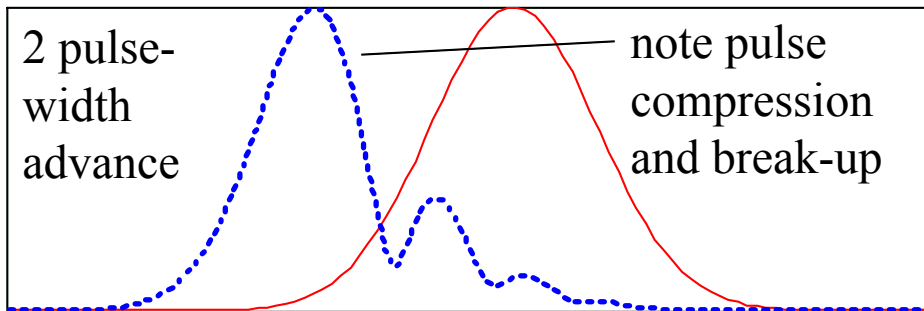
Slow light: single Lorentzian gain line

$$T = \exp(+16)$$

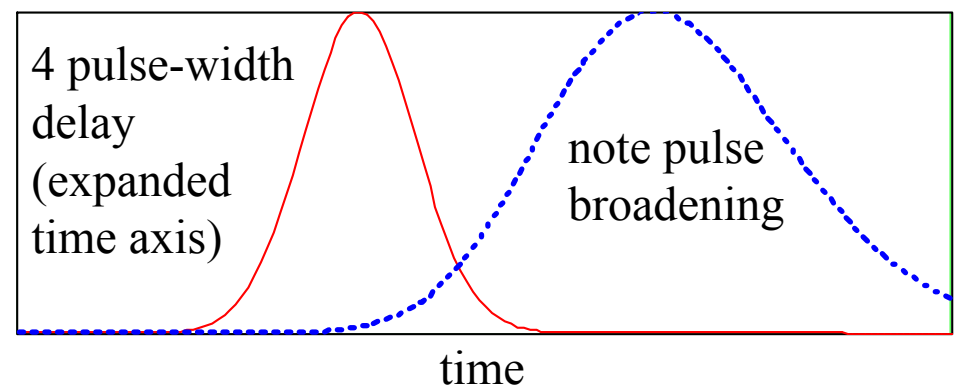
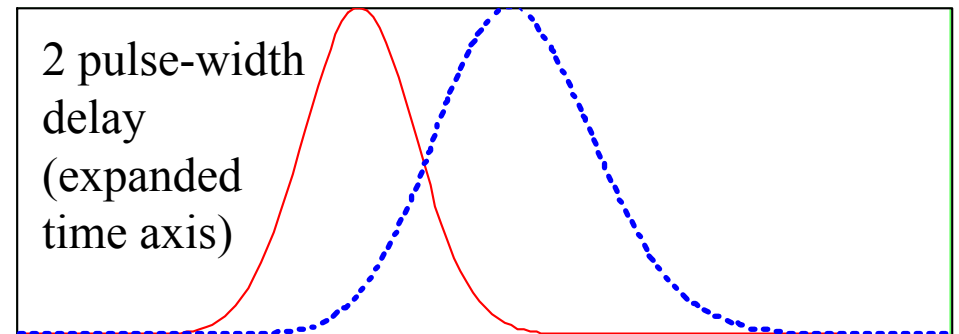
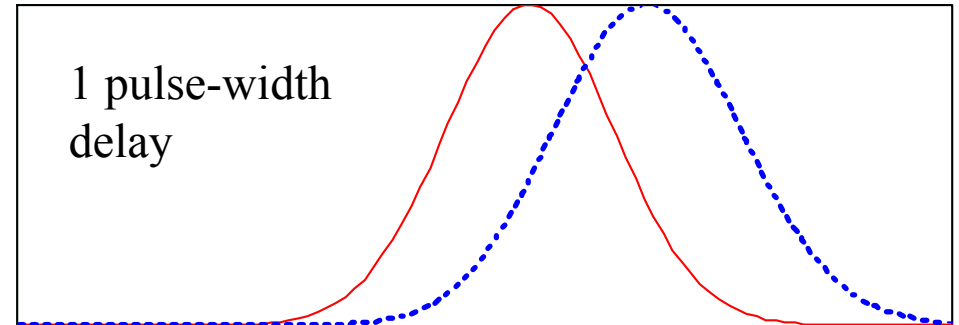
vary absorption line width to control delay

Same Gaussian input pulse in all cases

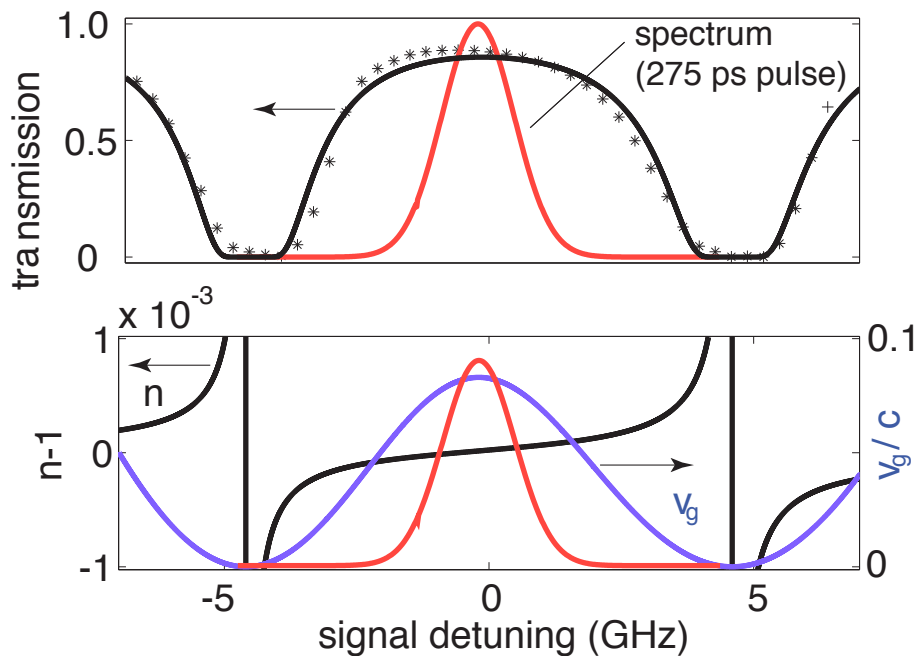
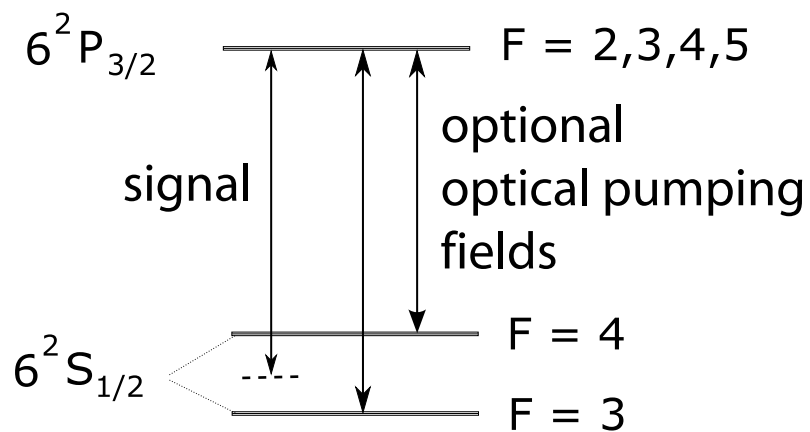
## Fast Light Results



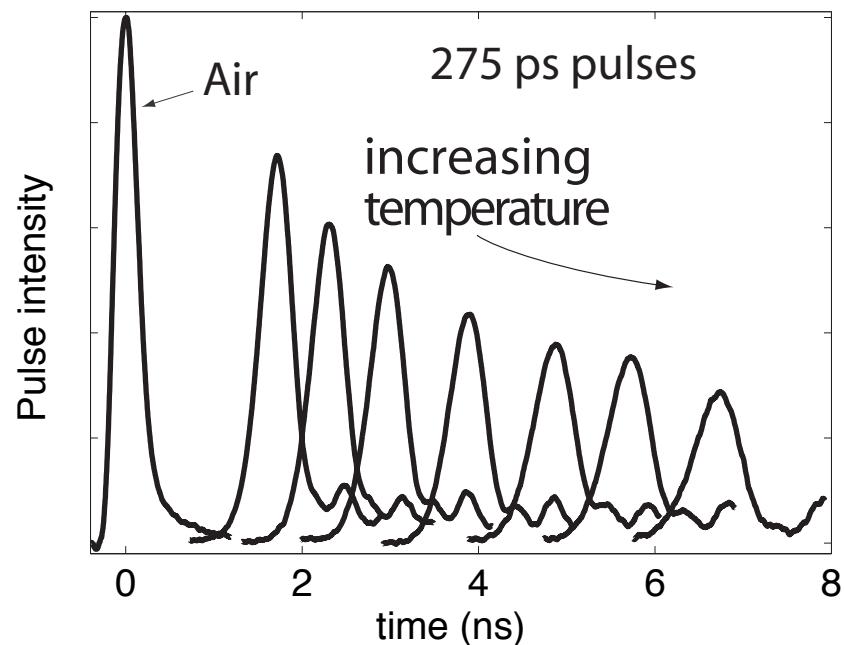
## Slow Light Results



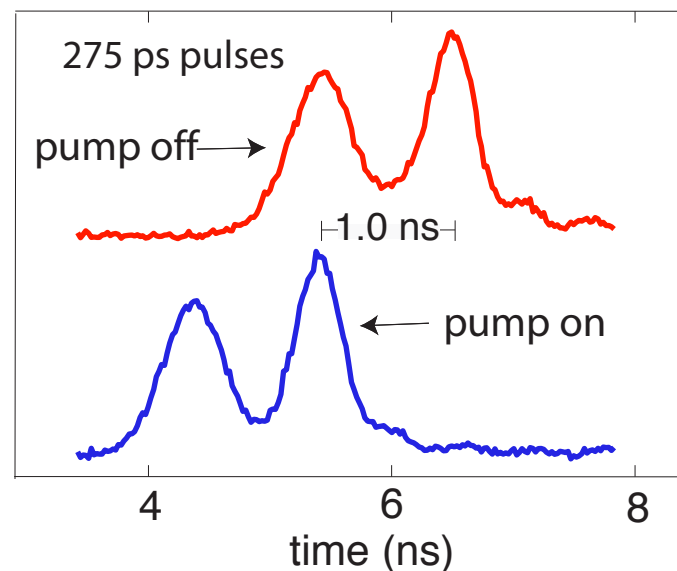
# Tunable Delays of up to 80 Pulse Widths in Atomic Cesium Vapor



- coarse tuning: temperature



- fine tuning: optical pumping

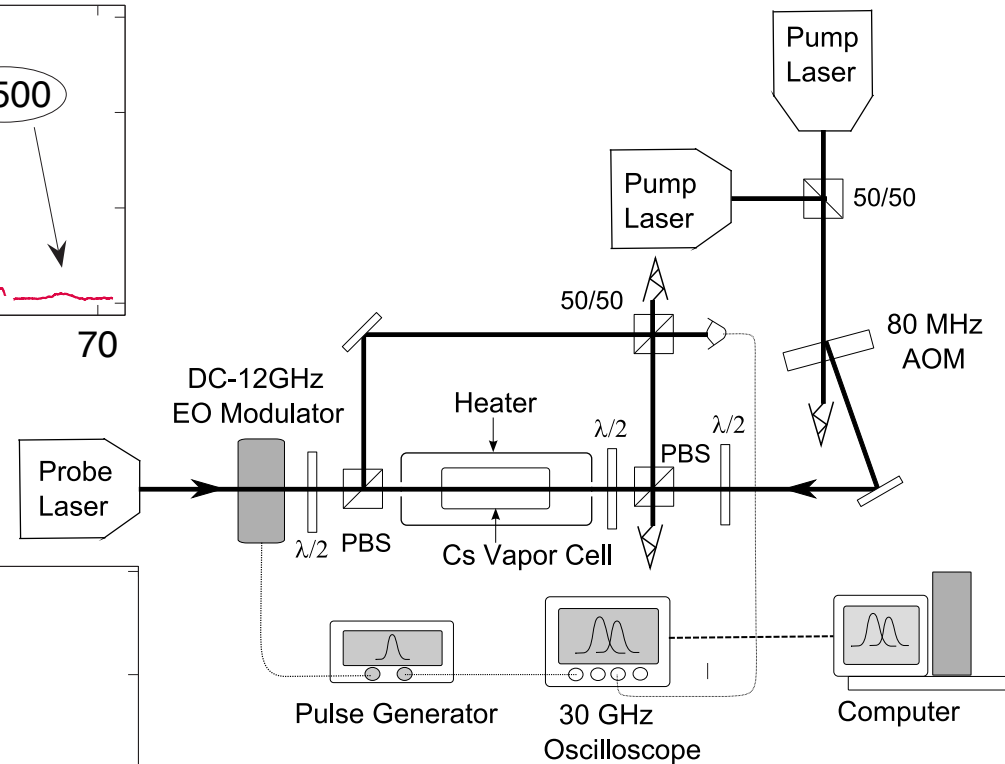
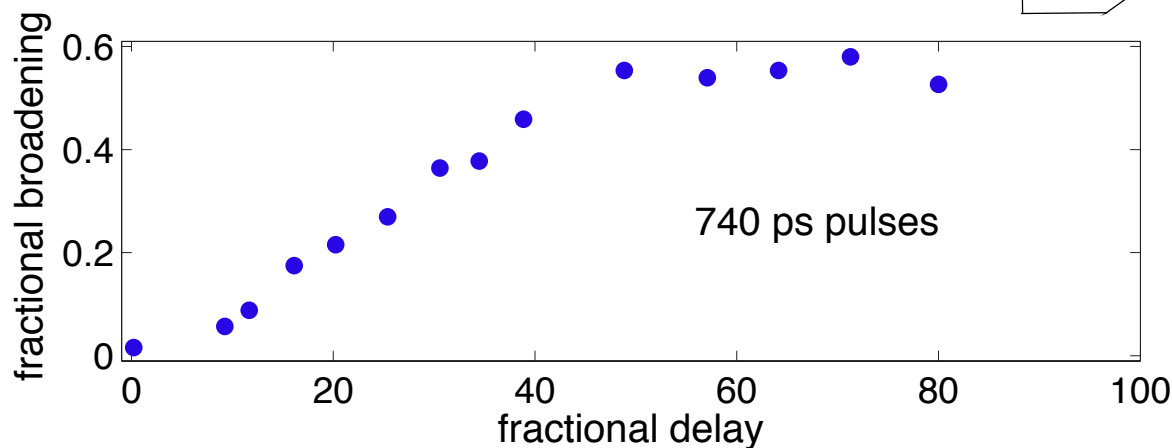
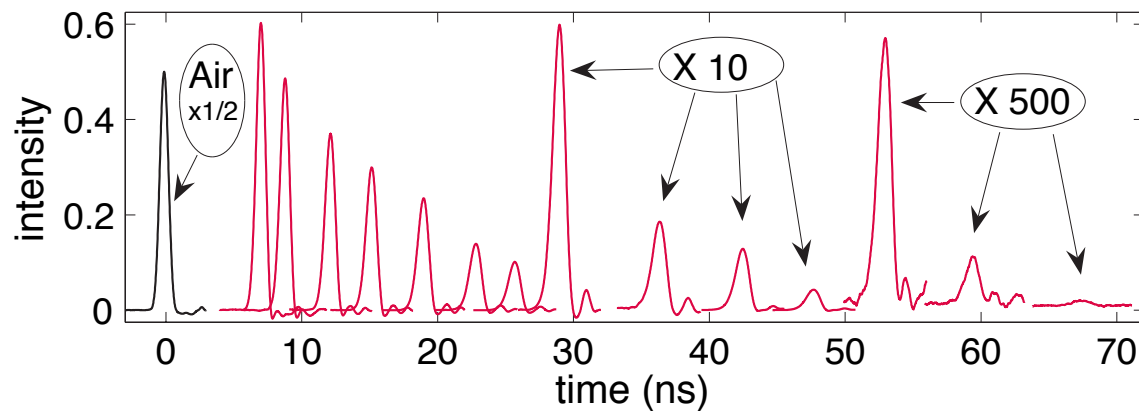


# Tunable Delays of up to 80 Pulse Widths in Atomic Cesium Vapor

Comment: In EIT based slow light, spectral reshaping is the dominant limitation. But far off resonance, this effect is negligible. Group velocity dispersion becomes important.

Longer input pulses lead to reduced gvd distortion and longer fractional delays

Results for 740 ps pulses





# How to Prevent Pulse Distortion (Which Can Limit Data Rates)

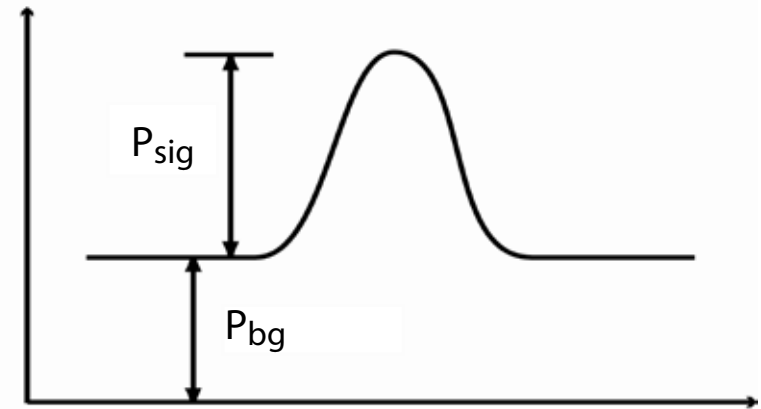
Two primary mechanisms for pulse distortion in EDFA

- Spectral broadening, leading to **temporal compression**  
CPO gain dip causes spectral components in the wings to be amplified more than central components

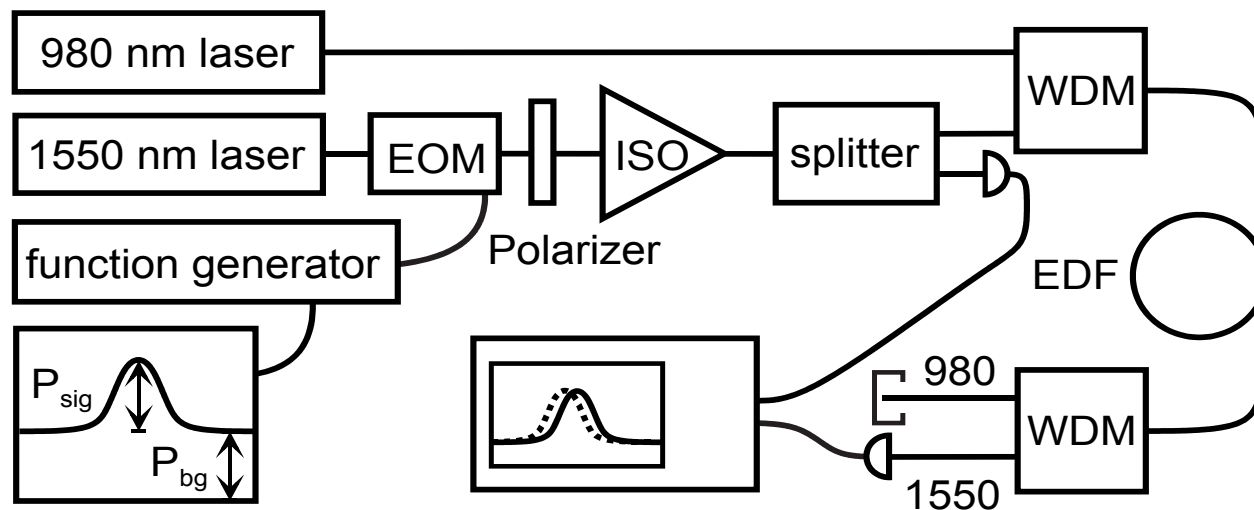
- Temporal gain recovery, leading to **temporal broadening**  
Leading edge of signal pulse saturates gain, but for long pulses, the trailing edge can experience recovered gain

To minimize second effect, add a cw background to reduce the influence of gain recovery

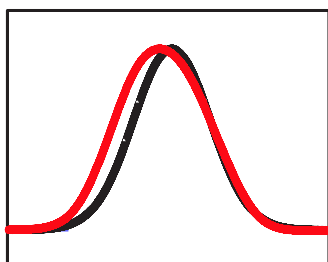
**For the proper choice of background power, the two effects exactly cancel!**



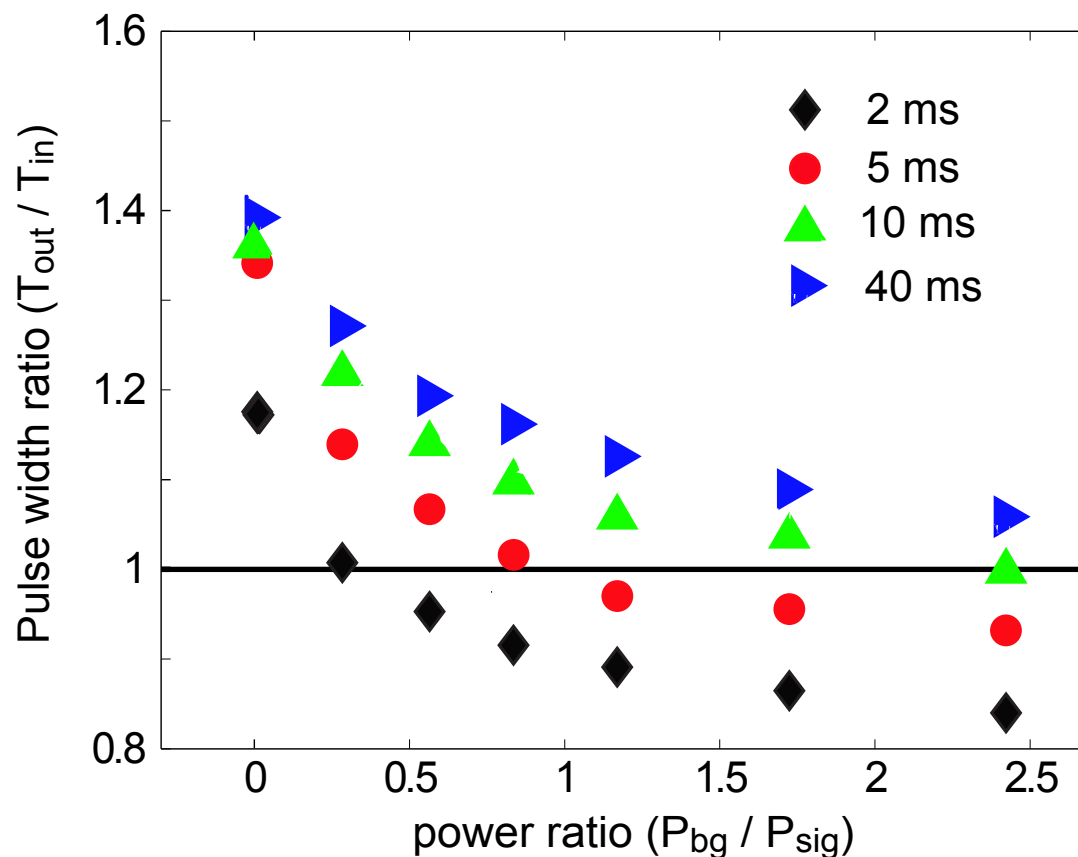
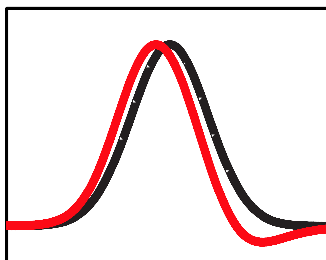
# Minimizing Pulse Distortion – Laboratory Results



broadening



compression



# Summary

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Slow-light techniques hold great promise for applications in telecommunications

Good progress being made in developing new slow-light techniques and applications

Backwards and superluminal propagation are strongly counterintuitive, but are fully explained by standard physics.

# Special Thanks to My Students and Research Associates

